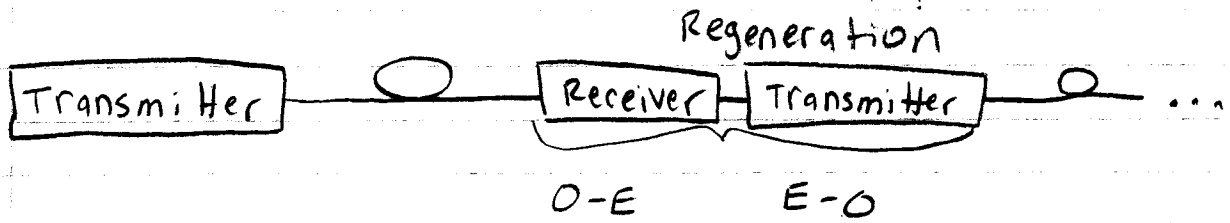


Chapter 17 Optical Amplification



O-E optical to electrical conversion
E-O electrical to optical conversion

What determines the distance between regeneration?

Signal-to-Noise Ratio : SNR

Dispersion

Signal power (attenuation)

What affects dispersion:

Fiber dispersion:

Dispersion reduction

Dispersion compensation

Laser linewidth: DFB lasers

Modulation: Direct modulation of lasers

External modulators

Signal power: Attenuation

Fiber

Component

OPTICAL AMPLIFIERS

Why optical amplifier vs. regeneration?

Format independent 1Gbps 40Gbps

Less complex

When will Optical Amplification not work?

Dispersion dominates

Too much noise

Optical amplifier amplifies noise

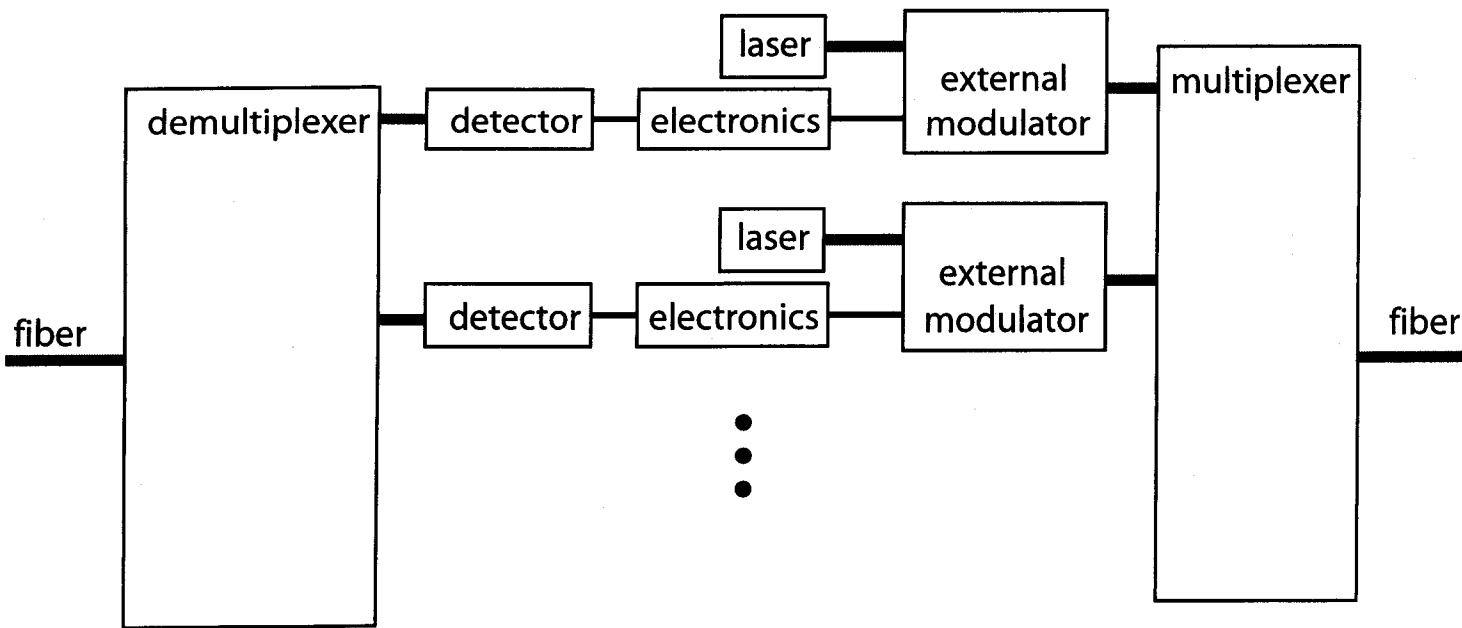
Optical amplifier generates noise

Criteria for Optical amplifier

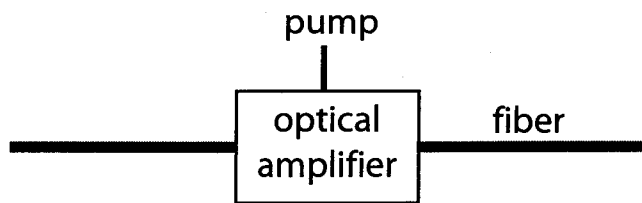
- 1. High power gain
- 2. High external pumping efficiency
- 3. Small saturation effect
- 4. Large bandwidth
- 5. Low added noise
- 6. Small crosstalk
- 7. Proper operating wavelength
- 8. Low insertion loss

Common amplifiers : Semiconductor amplifier
Erbium doped fiber amplifier

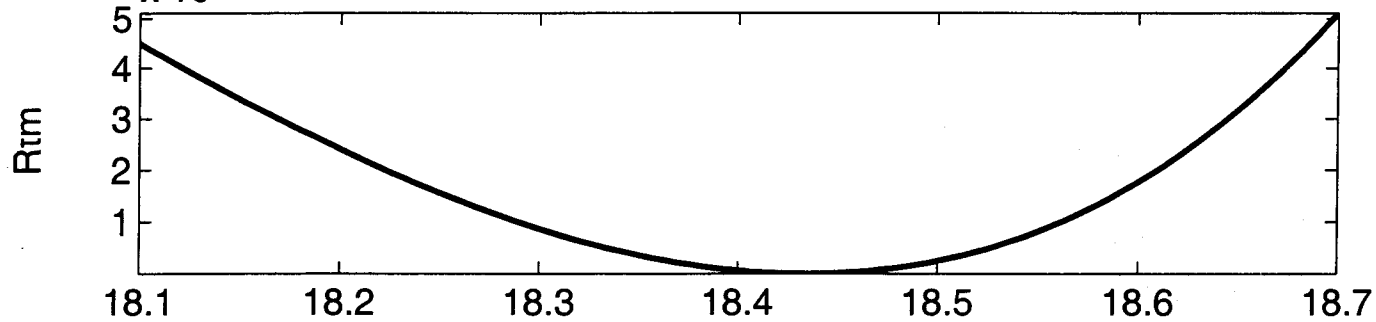
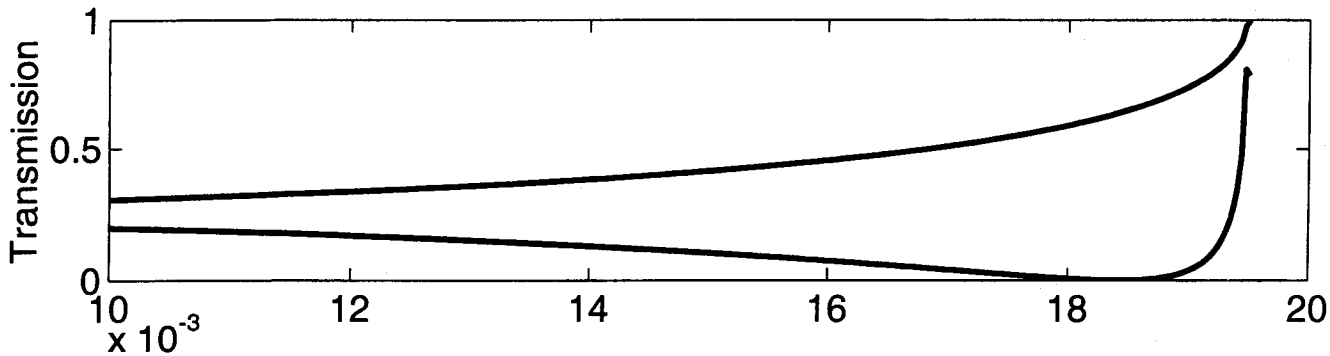
WDM Regenerator



Optical Amplifier

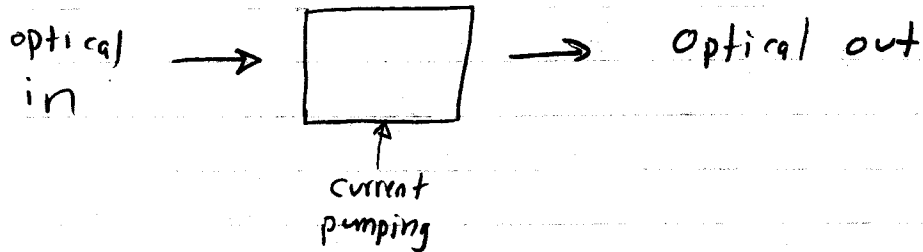


ni=3.0 nt=1.0



Semiconductor Optical Amplifier

Laser cavity biased below threshold



$$\text{gain (dB)} = \frac{(1 - r_L^2)(1 - r_R^2) G_0}{|1 - G_0 r_R r_L e^{\pm 2\beta L}|^2}$$

Variation in gain with wavelength

$$\text{variation} = \left(\frac{1 + G_0 r_R r_L}{1 - G_0 r_R r_L} \right)^2$$

Example:

For a gain of 20dB ($G_0 = 100$)
what is the minimum reflectivity for a
peak gain variation of 1dB

$$10^1 = 124 = \left(\frac{1 + 100 r_R r_L}{1 - 100 r_R r_L} \right)^2$$

$$\sqrt{124} (1 - 100 r_R r_L) = 1 + 100 r_R r_L$$

$$\frac{\sqrt{124} - 1}{100 + \sqrt{124} (100)} = r_R r_L = 5.37 \times 10^{-4}$$

Very small

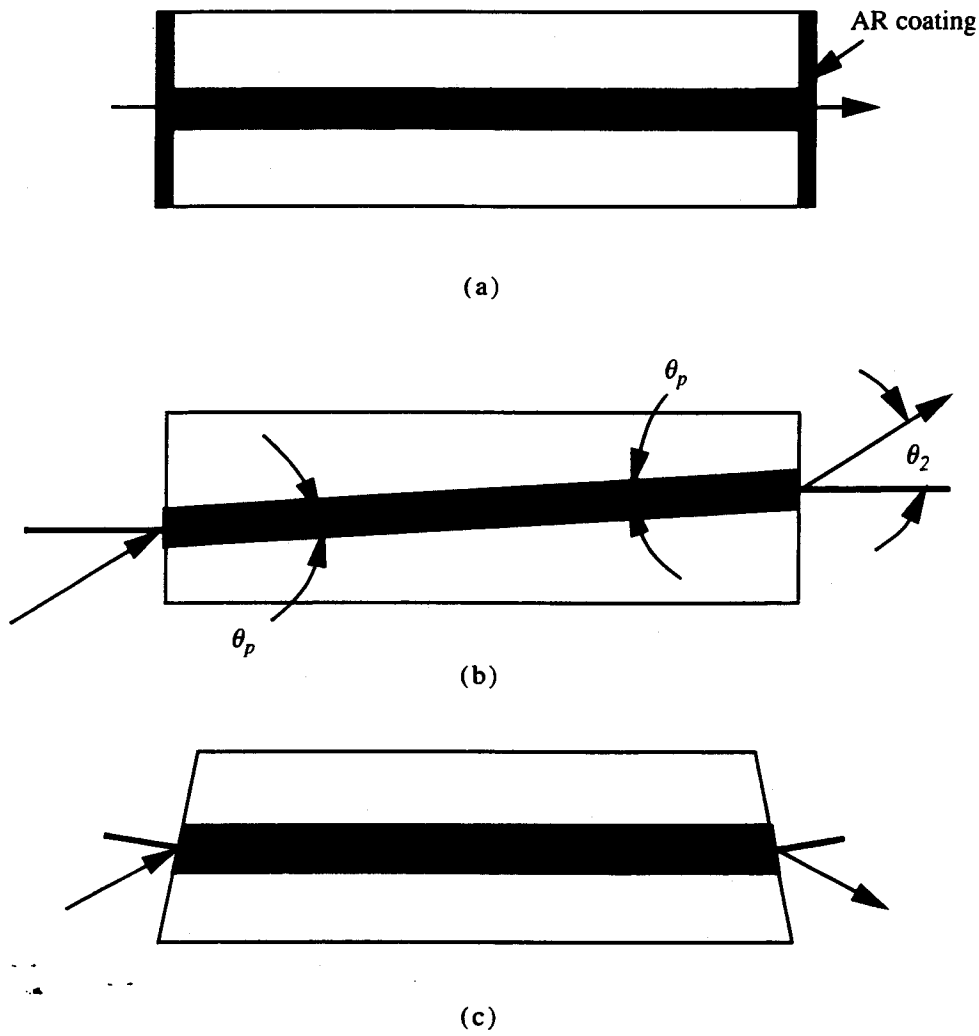
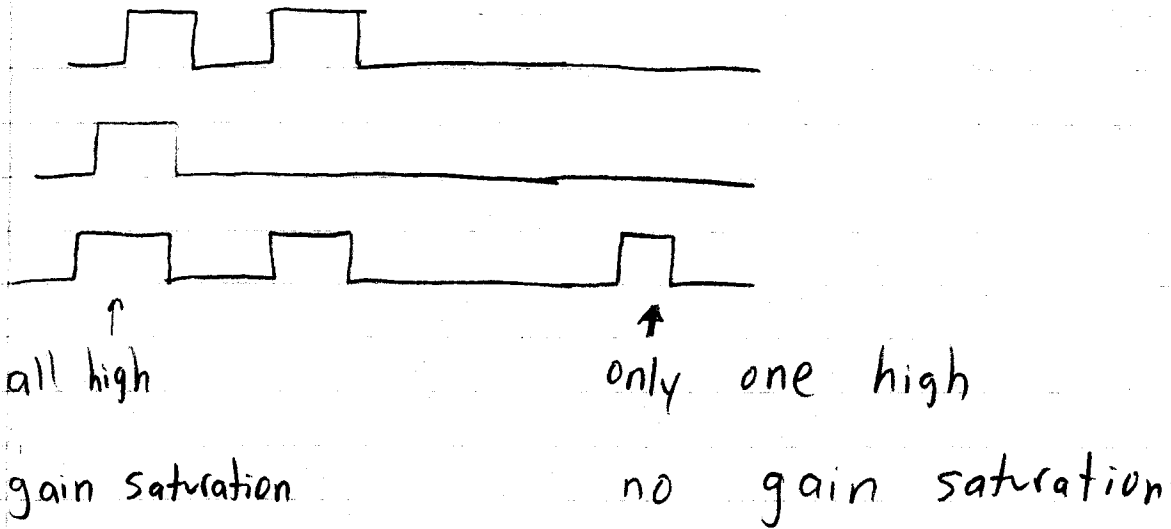


Figure 17.7 Methods to reduce the cavity end reflection: (a) antireflection coating, (b) cavity tilted at the Brewster angle, and (c) cavity facet tilted at the Brewster angle. In (b) and (c), note that the incident and output light is not parallel to the longitudinal direction of the cavity

With multiple wave length channels the carrier concentration varies randomly.



Variance of change in gain

$$E[\Delta g^2] \approx E[\Delta g_0^2] \quad T \gg \tau$$

$$E[\Delta g_0^2] \left(\frac{T}{\tau}\right)^3 \left(\frac{1}{3}\right) \quad T \ll \tau$$

$$\Delta g_0 = g(N) \frac{\Delta N_{ph}}{N_{ph,sat} + N_{ph,0}}$$

ΔN_{ph} : one channel
 $N_{ph,0}$: average

For semiconductor amplifiers

$$T \gg \tau$$

Semiconductor Optical Amplifiers (SOA)

High gain $\sim 20\text{dB}$
High Saturation power 5-10 dBm
Large BW
Can operate at any wavelength
Compact
Easy integration
Integration into arrays

Polarization dependent
High noise
cross talk between WDM channels

The crosstalk limits SOA in WDM applications

Erbium Doped Fiber Amplifier (EDFA)

Dope the core of an optical fiber with the rare earth ion Er^{+3}

The energy diagram has a meta-stable state with an emission wavelength around $\lambda = 1530 \text{ nm}$. and quick transitions into this state.

The population inversion is achieved by optically pumping at $\lambda = 980 \text{ nm}$ or 1480 nm .

The optical pumping is longitudinal rather than transverse resulting in better absorption of the optical pump.

The basic component of an EDFA are
pump laser
isolator

Coupler to combine signal and pump
Erbium doped fiber
Filter: to eliminate pump light

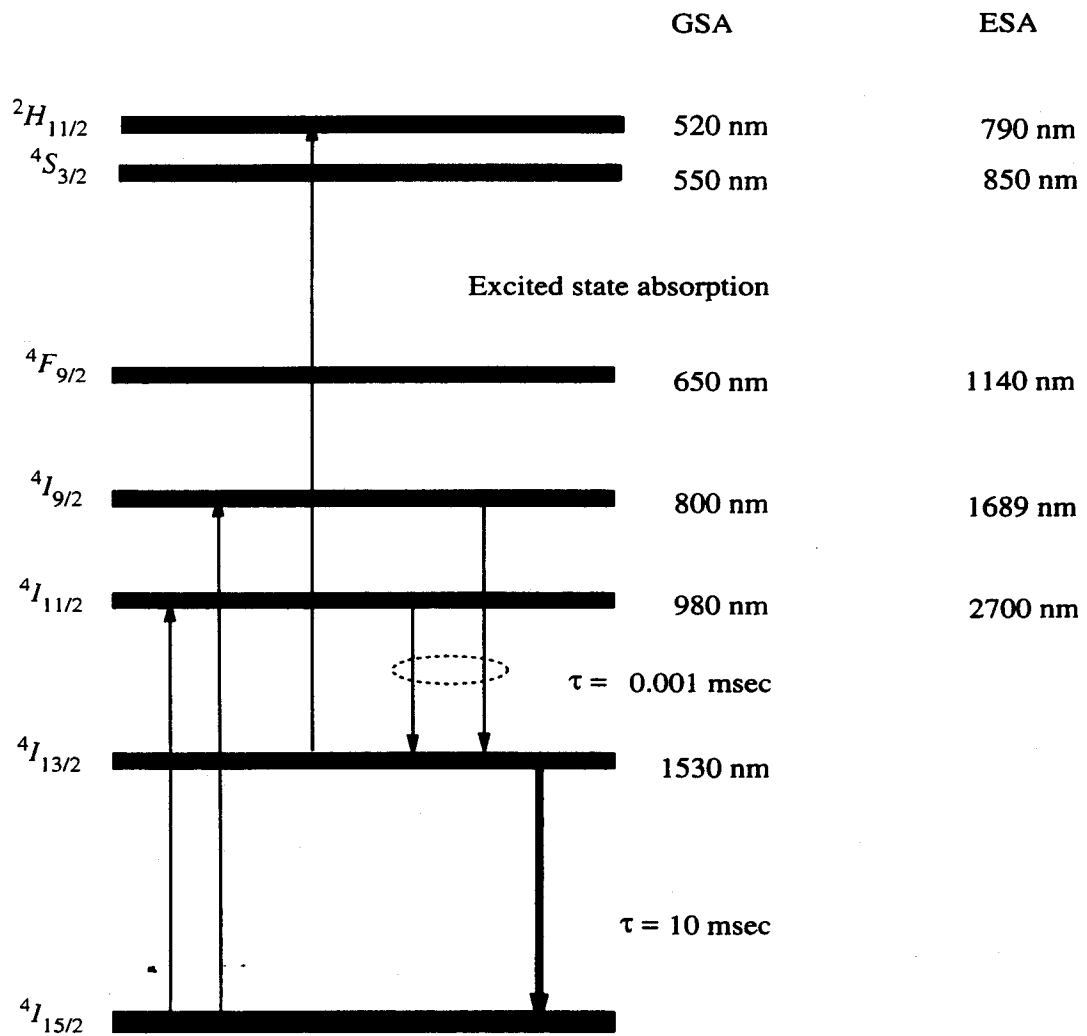
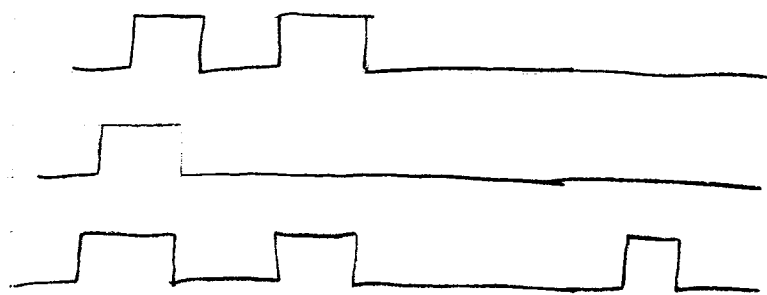


Figure 17.9 Energy levels of Er^{3+} . For each level, the ground state absorption (GSA) wavelength is the light wavelength needed to excite carriers from the ground state to the given level, and the excited state absorption (ESA) wavelength is the light wavelength needed to excite carriers from the given level to the metastable state $4I_{13/2}$.

With multiple wavelength channels the carrier concentration varies randomly.



↑
all high

↑
only one high

gain saturation

no gain saturation

Variance of change in gain

$$E[\Delta g^2] \approx E[\Delta g_0^2] \quad T \gg \tau$$

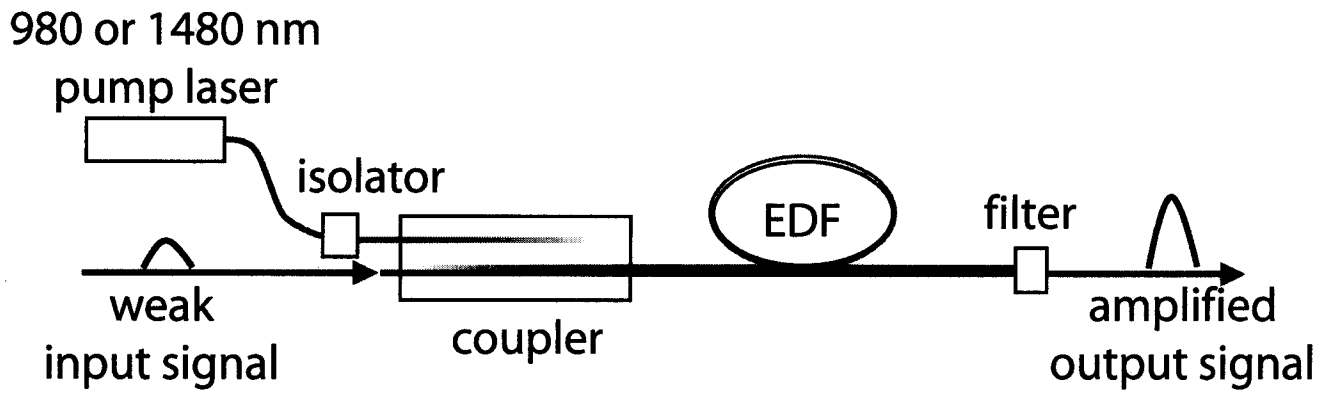
$$E[\Delta g_0^2] \left(\frac{\tau}{T}\right)^3 \left(\frac{1}{3}\right) \quad T \ll \tau$$

$$\Delta g_0 = g(N) \frac{\Delta N_{ph}}{N_{ph,sat} + N_{ph,0}}$$

ΔN_{ph} : one channel
 $N_{ph,0}$: average

For semiconductor amplifiers

$$T \gg \tau$$



The external pump rate is

$$R_p = W_p N_1$$

where N_1 is the ground state population

$$W_p = \frac{\sigma_a P_p}{h f_p A} \text{ s}^{-1}$$

$h f_p$ is the photon energy of the pump

A is the core area

σ_a is the absorption cross section

P_p is the pump power

The gain is given by

$$g = \sigma_e (N_2 - N_1)$$

N_2 is the concentration in the metastable state

N_1 is the concentration in the ground state

σ_e is the emission cross-section

The stimulated emission rate becomes

$$R_s = V_g g N_{ph}$$

$$P_{in} = (V_g)(N_{ph} A)(h f_s)$$

$$R_s = \frac{\sigma_e P_{in}}{h f_s A} (N_2 - N_1) \stackrel{\Delta}{=} W_s (N_2 - N_1)$$

$$W_s \stackrel{\Delta}{=} \frac{\sigma_e P_{in}}{h f_s A} \text{ s}^{-1}$$

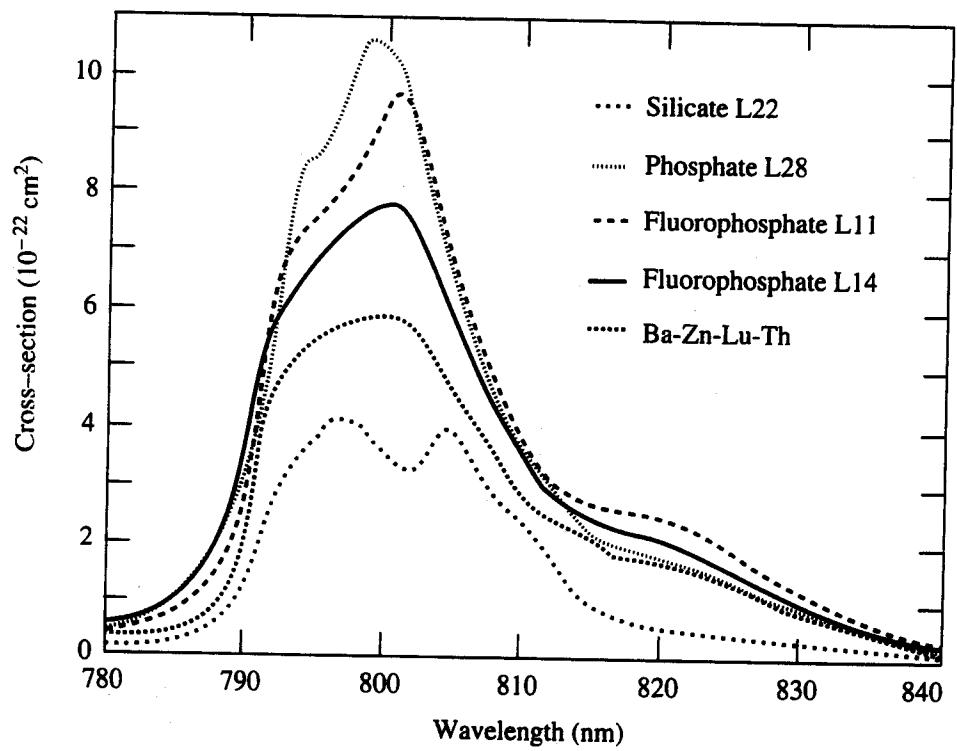


Figure 17.13 Absorption cross section of Er³⁺ at 800 nm.

Source: Reprinted, by permission, from Miniscalco, "Erbium-Doped Glasses for Fiber Amplifiers at 1500 nm," Figure 11 [10]. © 1991 by IEEE.

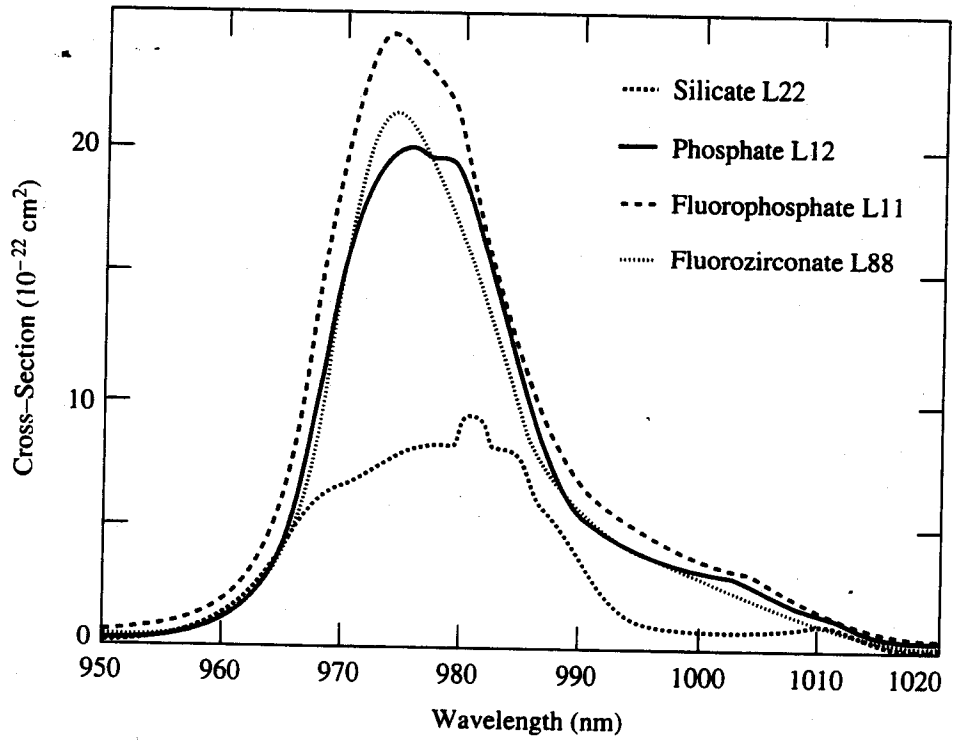


Figure 17.14 Absorption cross section of Er³⁺ at 980 nm.

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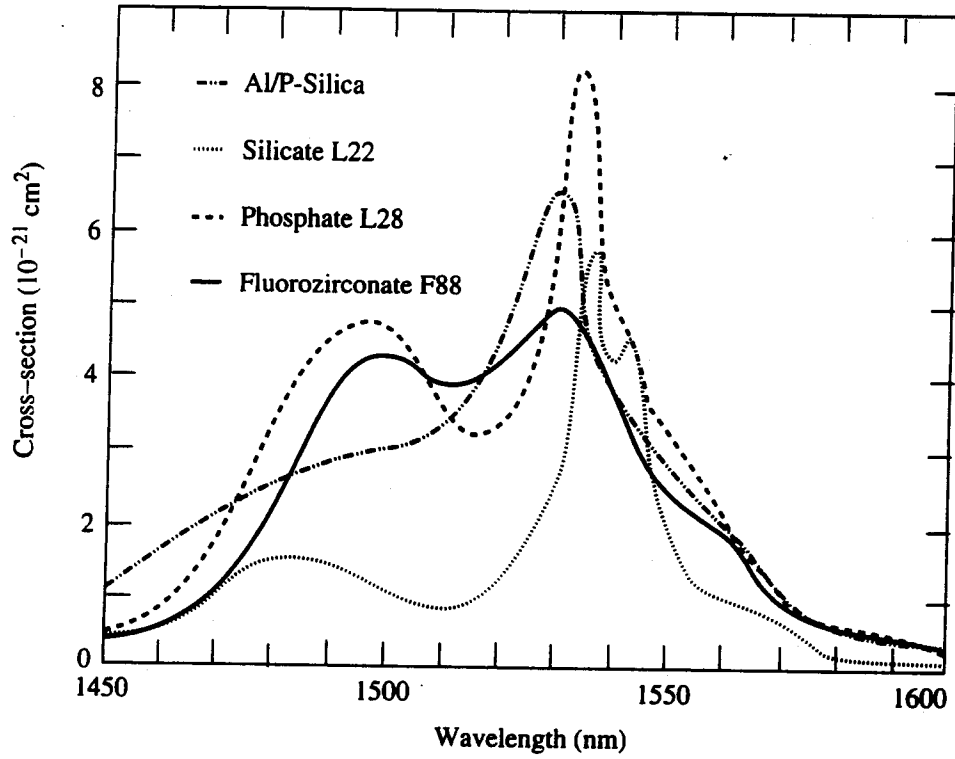


Figure 17.15 Absorption cross section of Er^{3+} at 1450 nm.

Source: Reprinted, by permission, from Miniscalco, "Erbium-Doped Glasses for Fiber Amplifiers at 1500 nm," Figure 7 [10]. © 1991 by IEEE.

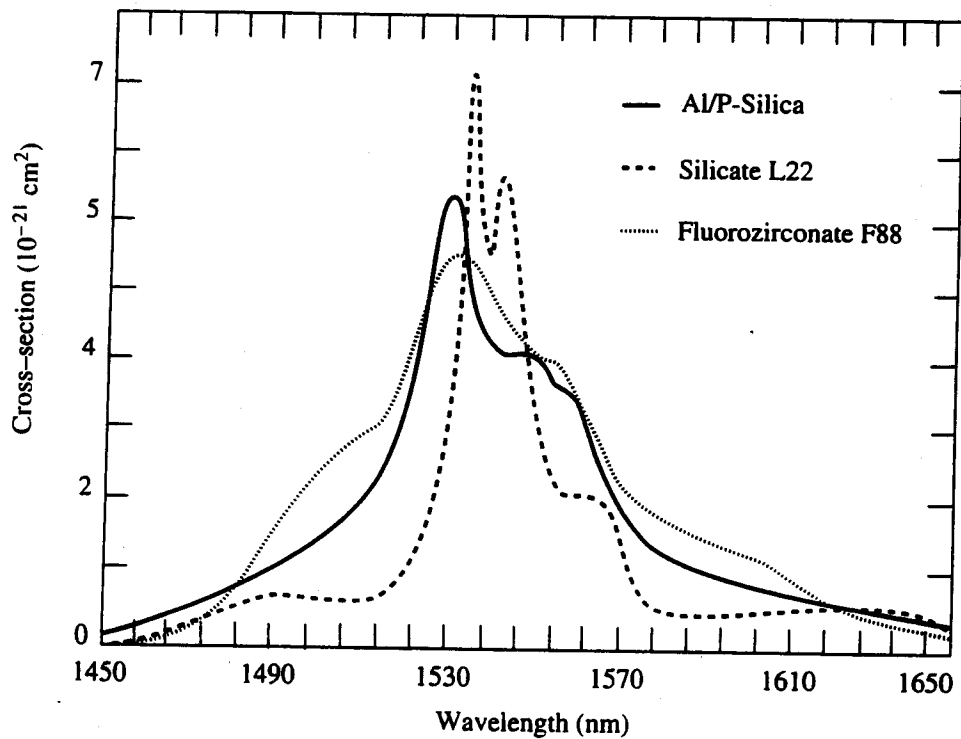


Figure 17.16 Emission cross section of Er^{3+} at 1540 nm.

Source: Reprinted, by permission, from Miniscalco, "Erbium-Doped Glasses for Fiber

The carrier rate equation is

$$\frac{dN_2}{dt} = \underset{\substack{\uparrow \\ \text{pump} \\ \text{rate}}}{W_p N_1} - \underset{\substack{\uparrow \\ \text{net stimulated} \\ \text{emission rat}}}{W_s (N_2 - N_1)} - \underset{\substack{\uparrow \\ \text{spontaneous recombination} \\ \text{rate}}}{\frac{N_2}{\tau_{sp}}} = - \frac{dN_1}{dt}$$

$$\tau_{sp} \sim 10 \text{ ns}$$

at steady state

$$0 = W_p N_1 - W_s (N_2 - N_1) - \frac{N_2}{\tau_{sp}}$$

$$0 = 2W_p N_1 - 2W_s (N_2 - N_1) - \frac{2N_2}{\tau_{sp}}$$

$$0 = W_p (N_1 + N_2) - W_p (N_2 - N_1) - 2W_s (N_2 - N_1) - \frac{1}{\tau_{sp}} (N_2 - N_1) - \frac{1}{\tau_{sp}} (N_1 + N_2)$$

$$(N_2 - N_1) (W_p + 2W_s + 1/\tau_{sp}) = (N_1 + N_2) (W_p - 1/\tau_{sp})$$

$$N_2 - N_1 = \frac{(W_p - 1/\tau_{sp})}{W_p + 2W_s + 1/\tau_{sp}} N_t$$

$N_t = N_1 + N_2$ is the total carrier density

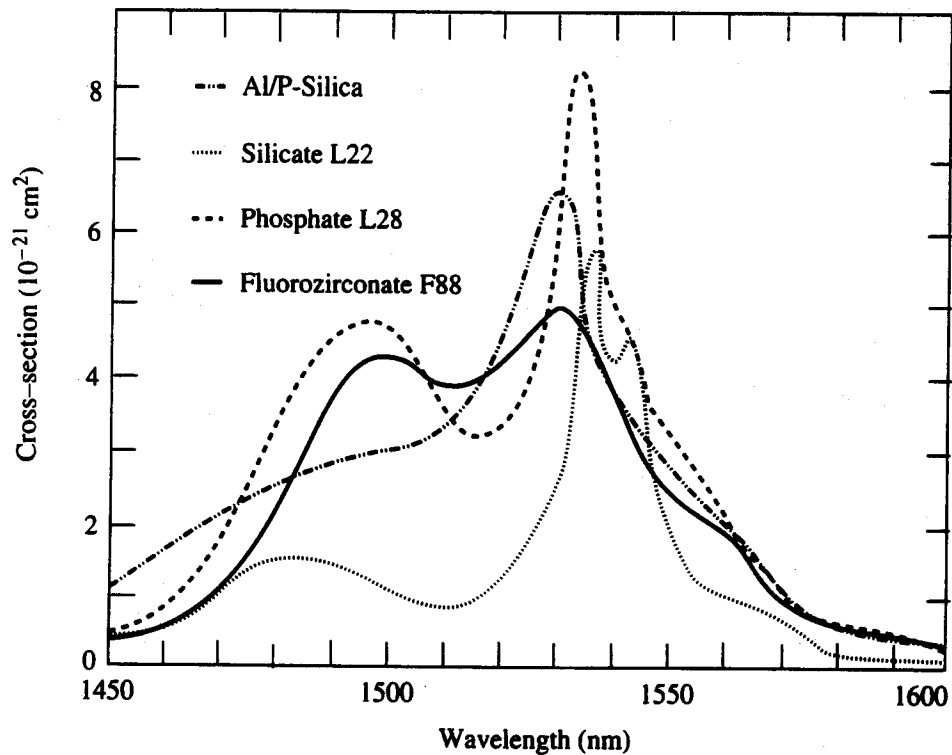


Figure 17.15 Absorption cross section of Er^{3+} at 1450 nm.

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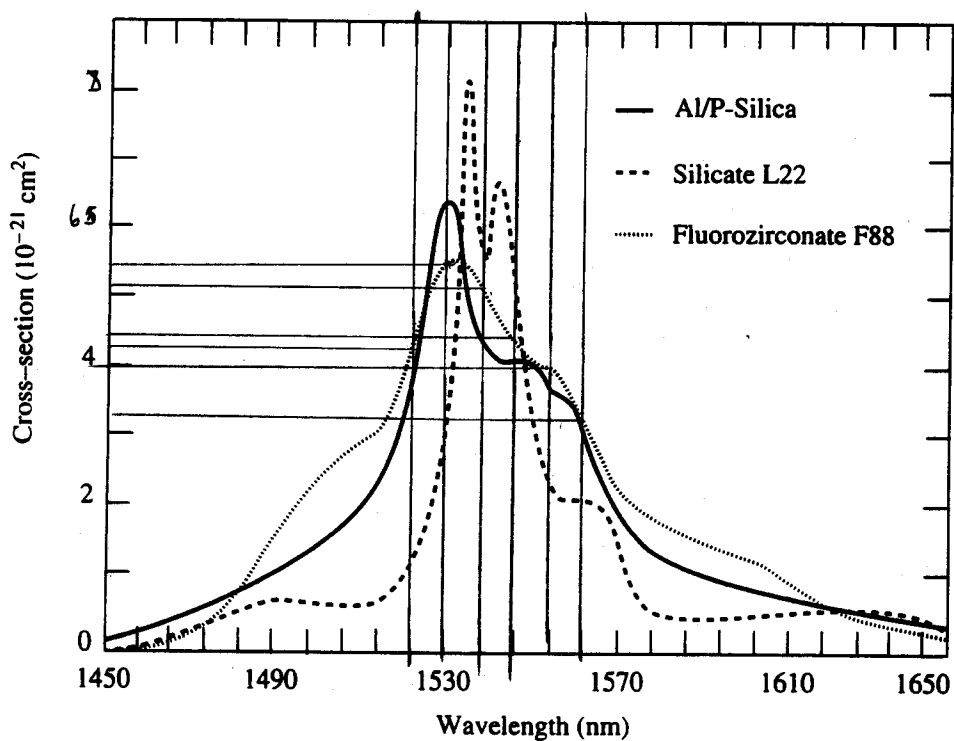


Figure 17.16 Emission cross section of Er^{3+} at 1540 nm.

Source: Reprinted, by permission, from Miniscalco, "Erbium-Doped Glasses for Fiber Amplifiers at 1500 nm," Figure 8 [10]. © 1991 by IEEE.

With a high pumping rate

$$W_p \gg W_s \quad W_p \gg \nu_{TSP}$$

$$N_2 - N_1 = \frac{W_p}{W_p} N_t$$

$$g = \sigma_e (N_2 - N_1) \approx \boxed{\sigma_e N_t \triangleq g^*}$$

$$g = \sigma_e (N_2 - N_1)$$

$$g = \sigma_e \left[\frac{W_p - \nu_{TSP}}{W_p + 2W_s + \nu_{TSP}} \right] N_t \quad N_t = \frac{g^*}{\sigma_e}$$

$$g = g^* \left[\frac{W_p - \nu_{TSP}}{W_p + 2W_s + \nu_{TSP}} \right]$$

Example :

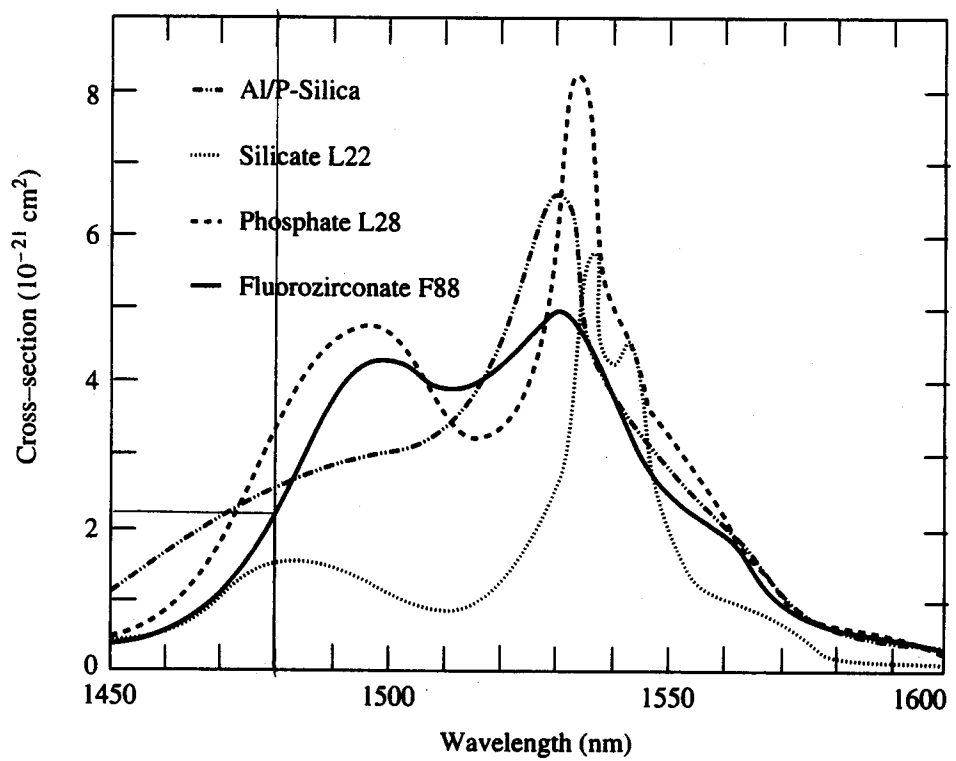


Figure 17.15 Absorption cross section of Er³⁺ at 1450 nm.

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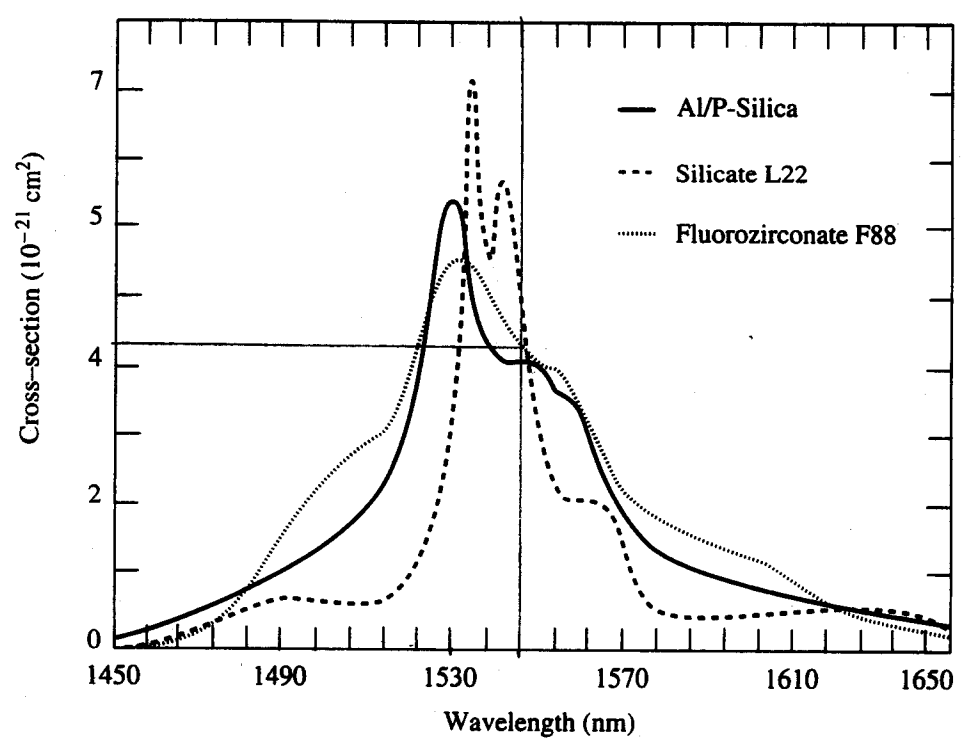


Figure 17.16 Emission cross section of Er³⁺ at 1540 nm.

Source: Reprinted, by permission, from Miniscalco, "Erbium-Doped Glasses for Fiber Amplifiers at 1500 nm," Figure 8 [10]. © 1991 by IEEE.

Pump at $\lambda_p = 1480 \text{ nm}$ $\sigma_a = 2.2 \times 10^{-21} \text{ cm}^2$

Signal at $\lambda = 1522, 1530, 1538, 1546, 1554, 1562 \text{ nm}$

$L = 5 \text{ m}$

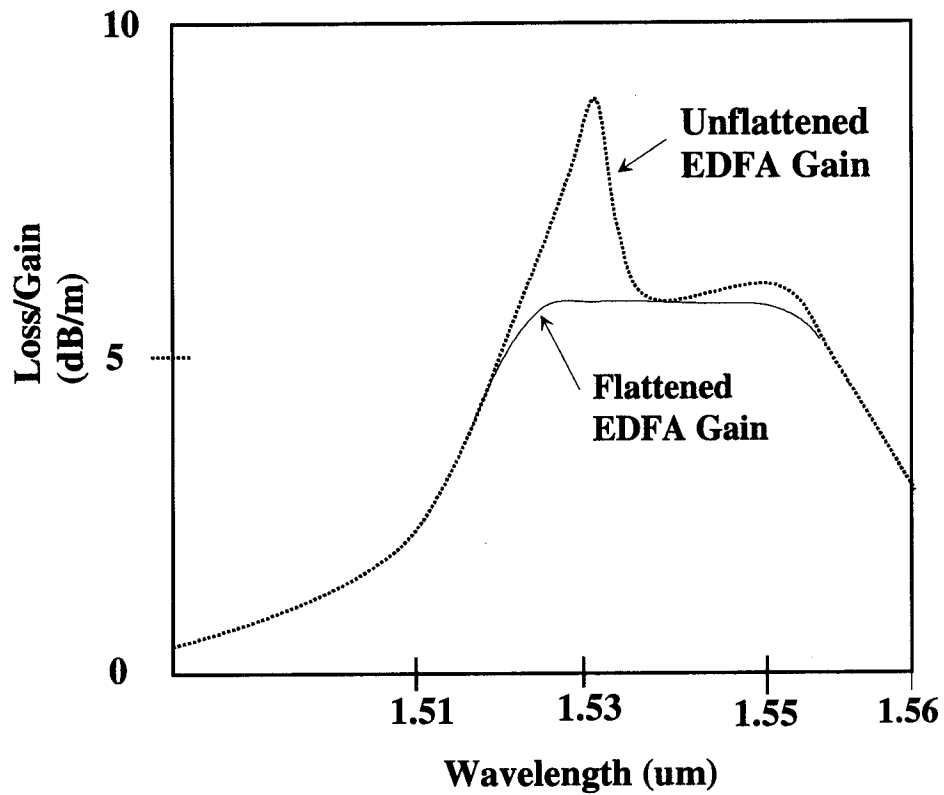
$\lambda_f \text{ (nm)}$	$\sigma_e \times 10^{-21} \text{ cm}^2$	$g^* \text{ (m}^{-1}\text{)}$	$g \text{ (m}^{-1}\text{)}$	$G \text{ (dB)}$
1522	4.28	3.42	1.77	38
1530	5.56	4.45	2.28	49
1538	5.22	4.18	2.14	46
1546	4.44	3.55	1.83	40
1554	4	3.20	1.65	36
1562	3.22	2.58	1.34	29

$$g^* = \sigma_e n_t$$

$$g = g^* \frac{W_p - \frac{1}{T_{sp}}}{W_p + 2W_s + \frac{1}{T_{sp}}}$$

$$G = (4.34 g) L$$

EDFA Gain Equalization



- Gain equalization can be accomplished in several ways:
 - a. Thin film filters
 - b. Long period fiber gratings
 - c. Chirped fiber Bragg gratings

The pump power is actually spatially dependent

$$\frac{dP_p}{dz} = -\sigma_a N_1 P_p$$

To relate N_1 to the total carrier density N_t
use Eq 17.33

$$N_2 - N_1 = \frac{(W_p - \gamma_{Tsp}) N_t}{W_p + 2W_s + \gamma_{Tsp}}$$

$$W_p(N_2 - N_1) + 2W_s(N_2 - N_1) + \gamma_{Tsp}(N_2 - N_1) = W_p(N_2 + N_1) - \gamma_{Tsp}(N_2 + N_1)$$

$$W_p(-2N_1) + 2W_s(N_2 - N_1) + \gamma_{Tsp}(2N_2) = 0$$

$$-N_1 W_p + W_s(N_2 + N_1 - 2N_1) + \gamma_{Tsp}(N_2 + N_1 - N_1) = 0$$

$$-N_1(W_p + 2W_s + \gamma_{Tsp}) + N_t(W_s + \gamma_{Tsp}) = 0$$

$$N_1 = \frac{(W_s + \gamma_{Tsp}) N_t}{W_p + 2W_s + \gamma_{Tsp}}$$

use this for pump power

$$\frac{dP_p}{dz} = -\sigma_a N_t P_p \frac{W_s + \gamma_{Tsp}}{W_p + 2W_s + \gamma_{Tsp}}$$

The signal power is also spatially varying

$$\frac{P_{in}}{dz} = g P_{in} \approx \sigma_e (N_2 - N_1) P_{in}$$

using Eq. 17.33

$$\frac{dP_{in}}{dz} = g^* \frac{W_p - \frac{1}{2} \Gamma_{sp}}{W_p + 2W_s + \frac{1}{2} \Gamma_{sp}} P_{in} = \frac{g_0}{1 + W_s/W_{sat}} P_{in}$$

$$g_0 = g^* \frac{W_p - \frac{1}{2} \Gamma_{sp}}{W_p + \frac{1}{2} \Gamma_{sp}}$$

$$W_{sat} = \frac{1}{2} (W_p + \frac{1}{2} \Gamma_{sp})$$

W_p gets smaller along the propagation direction
Gain saturation effect is stronger at the
output end of the amplifier

The optimum amplifier length is a function
of the pump power

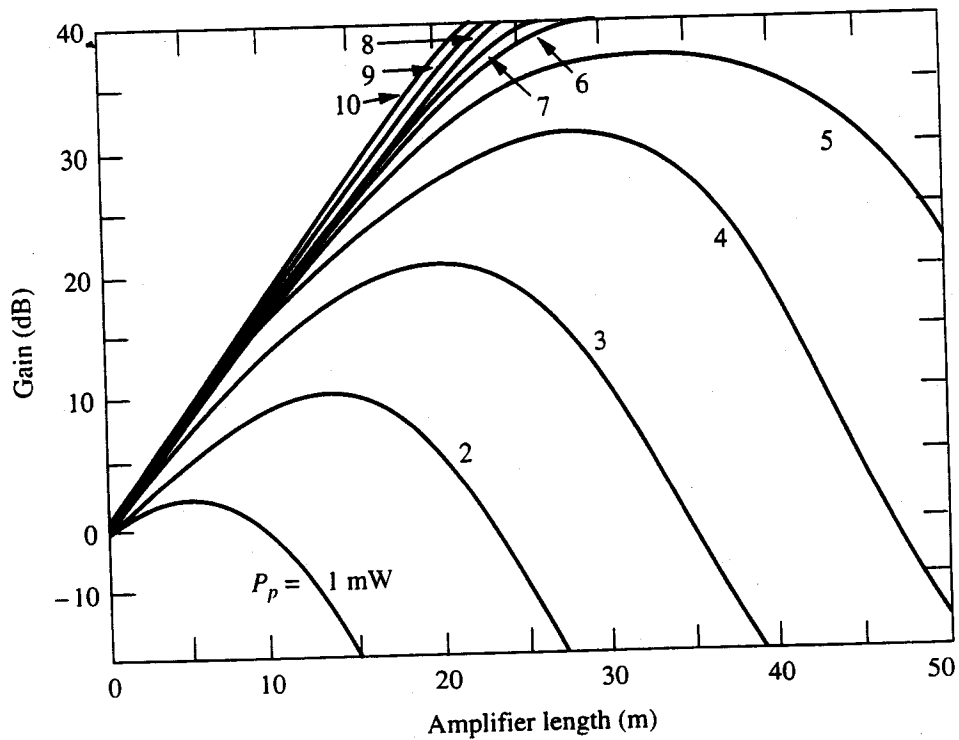


Figure 17.17 EDFA gain as a function of the amplifier length.

Source: Reprinted, by permission, from Giles and Desurvire, "Modeling Erbium-Doped Fiber Amplifiers," Figure 8b [12]. © 1991 by IEEE.

Erbium doped fiber amplifiers (EDFA)

High power transfer from pump to signal ($>50\%$)

Saturation output 10 - 25 dBm

Long gain constant ($\tau > 10\text{ms}$) to overcome patterning effect and crosstalk

Polarization independent

Already in fiber domain

Large devices m - Km

Amplified spontaneous emission (ASE) noise